Chebyshev Approximation of $(1 + 2x)\exp(x^2)$ erfc x in $0 \le x < \infty$

By M. M. Shepherd and J. G. Laframboise

Abstract. We have obtained a single Chebyshev expansion of the function $f(x) = (1 + 2x)\exp(x^2)\operatorname{erfc} x$ in $0 < x < \infty$, accurate to 22 decimal digits. The presence of the factors $(1 + 2x)\exp(x^2)$ causes f(x) to be of order unity throughout this range, ensuring that the use of f(x) for approximating $\operatorname{erfc} x$ will give uniform relative accuracy for all values of x.

I. Introduction. The functions $\operatorname{erfc} x = (2/\sqrt{\pi})\int_x^{\infty} \exp(-t^2) dt$ and $\exp(x^2)\operatorname{erfc} x$ occur frequently in kinetic theory of gases and related subjects. Calculation of these functions using the identity $\operatorname{erfc} x = 1 - \operatorname{erf} x$, together with available approximations [1] for $\operatorname{erf} x$, usually results in large relative errors for large x because $\operatorname{erf} x \to 1$ as $x \to \infty$. To overcome this difficulty, Clenshaw [2], Luke [3], [4], and Schonfelder [5] have presented Chebyshev approximations in which the range $0 \le x < \infty$ is split into two ranges $0 \le x \le c$ and $c \le x < \infty$, with $\operatorname{erf} x$ being Chebyshev-approximated in $0 \le x \le c$, and $x \exp(x^2)\operatorname{erfc} x$ being Chebyshevapproximated in $c \le x < \infty$. Clenshaw [2] uses c = 4, 33 terms for $x \le 4$ and 18 terms for x > 4, and obtains an accuracy of twenty decimal places (20D). Corresponding figures for Luke [3], [4] and Schonfelder [5] are c = 3, 25 and 22 terms, and 20D; and c = 2, 27 and 43 terms, and 30D. These authors use various transformations t(x) to map $c \le x < \infty$ into $-1 \le t < 1$. Use of the identity $\operatorname{erfc}(-x) = 2 - \operatorname{erfc} x$ eliminates the need to approximate erfc x for negative x.

Schonfelder [5] has also presented a single 43-term Chebyshev expansion of $\exp(x^2)$ erfc x for the entire interval $0 \le x < \infty$, using a relation of the form t = (x - k)/(x + k) to map this interval into $-1 \le t < 1$. Oldham [6] has presented a simple approximation of $\sqrt{\pi}x \exp(x^2)$ erfc x, having a maximum relative error of one part in 7000 and suitable for hand calculation.

Whenever a function to be Chebyshev-approximated has a zero within its interval of definition or at either end of it, such an approximation is likely to give large relative errors near such a zero because the usual procedures for calculating Chebyshev coefficients minimize maximum *absolute* error. Accordingly, it is advantageous to multiply erfc x by factors which yield a product of order unity for all x in $(0, \infty)$ and then to Chebyshev-approximate this product function, because one will then obtain good uniformity of relative as well as absolute error. Our chosen function, $f(x) = (1 + 2x)\exp(x^2)\operatorname{erfc} x$, satisfies this criterion. It has limiting values

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of 1 and $2/\sqrt{\pi} \approx 1.13$ at x = 0 and $x \to \infty$, respectively. In comparison, Schonfelder's function $\exp(x^2)$ erfc x approaches 0 as $x \to \infty$. Furthermore, our choice contains no irrational coefficients, in contrast with the more obvious choice $(1 + \sqrt{\pi}x)\exp(x^2)$ erfc x, which $\to 1$ at x = 0 and $x \to \infty$. A graph of f(x) appears in Figure 1.



FIGURE 1

Graph of the function F(t) defined by the relation $f(x) = (1 + 2x)\exp(x^2)\operatorname{erfc} x$ together with the mapping t = (x - 3.75)/(x + 3.75).

We have used the same transformation as that of Schonfelder [5], i.e. t = (x - k)/(x + k) with k = 3.75, to map $0 \le x < \infty$ into $-1 \le t < 1$. Tests of various k values for our f(x) yielded results similar to his, namely that this value gives near-optimum convergence of the resulting Chebyshev series over the precision range of greatest interest, i.e. 8D to 18D. Our calculations were done in IBM quadruple precision, which yields a machine precision of 34D.

II. Calculation of Chebyshev Coefficients. We have used the usual [7] form of an *m*th-order Chebyshev expansion. Thus, the Chebyshev polynomials $T_j(t)$ are given by

(1)
$$T_j(t) = \cos(j \arccos t); \quad j = 0, 1, 2, \ldots$$

The above-mentioned f(x) and transformation from x to t define a function F(t) which is expanded as follows:

(2)
$$F(t) = \sum_{j=0}^{m} c_j T_j(t),$$

where

(3)
$$c_{j} = \frac{\sum_{k=0}^{m} F(t_{k}) T_{j}(t_{k})}{\|T_{j}\|^{2}},$$

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(4)
$$t_k = \cos\left(\frac{2k+1}{m+1}\frac{\pi}{2}\right); \quad k = 0, 1, 2, \dots, m,$$

(5)
$$||T_0||^2 = m + 1;$$
 $||T_i||^2 = \frac{1}{2}(m + 1)$ for $i > 0.$

In order to calculate the required values of f(x), we note that the Taylor expansion

(6)
$$\operatorname{erfc} x = 1 - \frac{2}{\sqrt{\pi}} \left(t - \frac{t^3}{3} + \frac{t^5}{5 \cdot 2!} - \frac{t^7}{7 \cdot 3!} + \frac{t^9}{9 \cdot 4!} - \cdots \right)$$

can be rearranged [8] into the form

(7)
$$\exp(x^2)\operatorname{erfc} x = \exp(x^2) - \frac{2x}{\sqrt{\pi}} \sum_{n=0}^{\infty} \frac{2^n x^{2n}}{1 \cdot 3 \cdot 5 \cdot \cdot \cdot (2n+1)},$$

the use of which is less sensitive to roundoff errors.

The asymptotic expansion

(8)
$$\exp(x^2)\operatorname{erfc} x = \frac{1}{x\sqrt{\pi}} \left(1 - \frac{1}{2x^2} + \frac{1\cdot 3}{(2x^2)^2} - \frac{1\cdot 3\cdot 5}{(2x^2)^3} + \cdots \right)$$

is of limited use when x is large. The continued-fraction expansion

(9)
$$\sqrt{\pi}x \exp(x^2) \operatorname{erfc} x = \frac{1}{|1|} + \frac{1}{|2x^2|} + \frac{2}{|1|} + \frac{2}{|1|} + \frac{3}{|2x^2|} + \cdots$$

(Perron [9]) yields better precision. Perron [10] gives the following algorithm for use of (9).

We define

(10)

$$A_{-1} = 1, \quad A_0 = b_0, \\
B_{-1} = 0, \quad B_0 = 1; \\
a_i = i/(2x^2), \quad b_i = 1 \quad \text{for } i = 0, 1, 2, 3, \dots; \\
b_0 + \frac{a_{1}}{|b_1|} + \frac{a_2}{|b_2|} + \dots + \frac{a_n}{|b_n|} = \frac{A_n}{B_n}.$$

Then A_n and B_n are given recursively by the relations

(11)
$$\begin{array}{c} A_n = b_n A_{n-1} + a_n A_{n-2} \\ B_n = b_n B_{n-1} + a_n B_{n-2} \end{array} \right\}, \qquad n = 1, 2, 3, \ldots$$

At smaller values of x, the convergence of (9) becomes slower. To overcome this, we have used double Aitken δ^2 extrapolation (Burden et al. [11, pp. 56-57]) as follows. If y_j is the approximation obtained by taking j terms of (9), then the sequences of numbers

(12)
$$y'_{j} = y_{j} - (y_{j} - y_{j-1})^{2} / (y_{j} - 2y_{j-1} + y_{j-2}),$$
$$y''_{j} = y'_{j} - (y'_{j} - y'_{j-1})^{2} / (y'_{j} - 2y'_{j-1} + y'_{j-2}),$$

converge progressively faster to (9). Use of (12) with (9)–(11) improved the fit of the Chebyshev approximation by about five orders of magnitude.

We have used (7) for $x \le 2.83$ and (9)–(12) for x > 2.83. This yielded values of f(x) accurate to at least 23D for all x. The resulting Chebyshev coefficients,

generated by (3)-(5), are shown in Table 1. Use of these in (2) gives an approximation which reproduces f(x) to at least 22D for all x. Schonfelder [5] generates Chebyshev coefficients using a different method [12], [13] in which an expression equivalent to (2) is substituted into a linear differential equation satisfied by the given function. Together with the boundary conditions satisfied by the same function, this procedure generates an infinite set of simultaneous linear equations for the c_i , a truncated version of which is then solved.

TABLE 1

Y=(1+2X) *EXP(X*X)*ERFC(X) T=(X-K)/(X+K) K= 3.75 X=(0, INF) ORD C(N) 0.1177578934567401754080Q+ -0.45900545806464773310 0.842491333665179155840 01 0.592099399981918904 98Q 2665866843530 75 0.9074997670705265094 241316354041 76081 6916973 025012 218864 010 .1076 002 631808834 286879501 .1945586 n 85 .965469675 334781 -18645 0 12507 0 506 â 20

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Department of Computer Science York University Toronto, Ontario, Canada M3J 1P3

Department of Physics York University Toronto, Ontario, Canada M3J 1P3

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