# Chebyshev Approximation of $(1+2 x) \exp \left(x^{2}\right) \operatorname{erfc} x$ <br> in $0 \leqslant x<\infty$ 

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#### Abstract

We have obtained a single Chebyshev expansion of the function $f(x)=$ $(1+2 x) \exp \left(x^{2}\right) \operatorname{erfc} x$ in $0 \leqslant x<\infty$, accurate to 22 decimal digits. The presence of the factors $(1+2 x) \exp \left(x^{2}\right)$ causes $f(x)$ to be of order unity throughout this range, ensuring that the use of $f(x)$ for approximating erfc $x$ will give uniform relative accuracy for all values of $x$.


I. Introduction. The functions erfc $x=(2 / \sqrt{\pi}) \int_{x}^{\infty} \exp \left(-t^{2}\right) d t$ and $\exp \left(x^{2}\right) \operatorname{erfc} x$ occur frequently in kinetic theory of gases and related subjects. Calculation of these functions using the identity erfc $x=1-\operatorname{erf} x$, together with available approximations [1] for erf $x$, usually results in large relative errors for large $x$ because erf $x \rightarrow 1$ as $x \rightarrow \infty$. To overcome this difficulty, Clenshaw [2], Luke [3], [4], and Schonfelder [5] have presented Chebyshev approximations in which the range $0 \leqslant x<\infty$ is split into two ranges $0 \leqslant x \leqslant c$ and $c \leqslant x<\infty$, with erf $x$ being Chebyshev-approximated in $0 \leqslant x \leqslant c$, and $x \exp \left(x^{2}\right) \operatorname{erfc} x$ being Chebyshevapproximated in $c \leqslant x<\infty$. Clenshaw [2] uses $c=4$, 33 terms for $x \leqslant 4$ and 18 terms for $x>4$, and obtains an accuracy of twenty decimal places (20D). Corresponding figures for Luke [3], [4] and Schonfelder [5] are $c=3,25$ and 22 terms, and 20D; and $c=2,27$ and 43 terms, and 30D. These authors use various transformations $t(x)$ to map $c \leqslant x<\infty$ into $-1 \leqslant t<1$. Use of the identity $\operatorname{erfc}(-x)=2-\operatorname{erfc} x$ eliminates the need to approximate erfc $x$ for negative $x$.

Schonfelder [5] has also presented a single 43-term Chebyshev expansion of $\exp \left(x^{2}\right) \operatorname{erfc} x$ for the entire interval $0 \leqslant x<\infty$, using a relation of the form $t=(x-k) /(x+k)$ to map this interval into $-1 \leqslant t<1$. Oldham [6] has presented a simple approximation of $\sqrt{\pi} x \exp \left(x^{2}\right) \operatorname{erfc} x$, having a maximum relative error of one part in 7000 and suitable for hand calculation.

Whenever a function to be Chebyshev-approximated has a zero within its interval of definition or at either end of it, such an approximation is likely to give large relative errors near such a zero because the usual procedures for calculating Chebyshev coefficients minimize maximum absolute error. Accordingly, it is advantageous to multiply erfc $x$ by factors which yield a product of order unity for all $x$ in $(0, \infty)$ and then to Chebyshev-approximate this product function, because one will then obtain good uniformity of relative as well as absolute error. Our chosen function, $f(x)=(1+2 x) \exp \left(x^{2}\right) \operatorname{erfc} x$, satisfies this criterion. It has limiting values

[^0]of 1 and $2 / \sqrt{\pi} \approx 1.13$ at $x=0$ and $x \rightarrow \infty$, respectively. In comparison, Schonfelder's function $\exp \left(x^{2}\right) \operatorname{erfc} x$ approaches 0 as $x \rightarrow \infty$. Furthermore, our choice contains no irrational coefficients, in contrast with the more obvious choice $(1+\sqrt{\pi} x) \exp \left(x^{2}\right) \operatorname{erfc} x$, which $\rightarrow 1$ at $x=0$ and $x \rightarrow \infty$. A graph of $f(x)$ appears in Figure 1.


Figure 1
Graph of the function $F(t)$ defined by the relation $f(x)=(1+2 x) \exp \left(x^{2}\right) \operatorname{erfc} x$ together with the mapping $t=(x-3.75) /(x+3.75)$.

We have used the same transformation as that of Schonfelder [5], i.e. $t=$ $(x-k) /(x+k)$ with $k=3.75$, to map $0 \leqslant x<\infty$ into $-1 \leqslant t<1$. Tests of various $k$ values for our $f(x)$ yielded results similar to his, namely that this value gives near-optimum convergence of the resulting Chebyshev series over the precision range of greatest interest, i.e. 8D to 18D. Our calculations were done in IBM quadruple precision, which yields a machine precision of 34D.
II. Calculation of Chebyshev Coefficients. We have used the usual [7] form of an $m$ th-order Chebyshev expansion. Thus, the Chebyshev polynomials $T_{j}(t)$ are given by

$$
\begin{equation*}
T_{j}(t)=\cos (j \operatorname{arc} \cos t) ; \quad j=0,1,2, \ldots \tag{1}
\end{equation*}
$$

The above-mentioned $f(x)$ and transformation from $x$ to $t$ define a function $F(t)$ which is expanded as follows:

$$
\begin{equation*}
F(t)=\sum_{j=0}^{m} c_{j} T_{j}(t) \tag{2}
\end{equation*}
$$

where

$$
\begin{equation*}
c_{j}=\frac{\sum_{k=0}^{m} F\left(t_{k}\right) T_{j}\left(t_{k}\right)}{\left\|T_{j}\right\|^{2}}, \tag{3}
\end{equation*}
$$

$$
\begin{gather*}
t_{k}=\cos \left(\frac{2 k+1}{m+1} \frac{\pi}{2}\right) ; \quad k=0,1,2, \ldots, m  \tag{4}\\
\left\|T_{0}\right\|^{2}=m+1 ; \quad\left\|T_{i}\right\|^{2}=\frac{1}{2}(m+1) \quad \text { for } i>0 \tag{5}
\end{gather*}
$$

In order to calculate the required values of $f(x)$, we note that the Taylor expansion

$$
\begin{equation*}
\operatorname{erfc} x=1-\frac{2}{\sqrt{\pi}}\left(t-\frac{t^{3}}{3}+\frac{t^{5}}{5 \cdot 2!}-\frac{t^{7}}{7 \cdot 3!}+\frac{t^{9}}{9 \cdot 4!}-\cdots\right) \tag{6}
\end{equation*}
$$

can be rearranged [8] into the form

$$
\begin{equation*}
\exp \left(x^{2}\right) \operatorname{erfc} x=\exp \left(x^{2}\right)-\frac{2 x}{\sqrt{\pi}} \sum_{n=0}^{\infty} \frac{2^{n} x^{2 n}}{1 \cdot 3 \cdot 5 \cdots(2 n+1)} \tag{7}
\end{equation*}
$$

the use of which is less sensitive to roundoff errors.
The asymptotic expansion

$$
\begin{equation*}
\exp \left(x^{2}\right) \operatorname{erfc} x=\frac{1}{x \sqrt{\pi}}\left(1-\frac{1}{2 x^{2}}+\frac{1 \cdot 3}{\left(2 x^{2}\right)^{2}}-\frac{1 \cdot 3 \cdot 5}{\left(2 x^{2}\right)^{3}}+\cdots\right) \tag{8}
\end{equation*}
$$

is of limited use when $x$ is large. The continued-fraction expansion

$$
\begin{equation*}
\sqrt{\pi} x \exp \left(x^{2}\right) \operatorname{erfc} x=\frac{1}{\sqrt{1}}+\frac{\left.\frac{1}{2 x^{2}} \right\rvert\,}{\mid 1}+\frac{\left.\frac{2}{2 x^{2}} \right\rvert\,}{\mid 1}+\frac{\frac{3}{2 x^{2}}}{\sqrt{\mid}}+\cdots \tag{9}
\end{equation*}
$$

(Perron [9]) yields better precision. Perron [10] gives the following algorithm for use of (9).

We define

$$
\begin{align*}
& A_{-1}=1, \quad A_{0}=b_{0}, \\
& B_{-1}=0, \quad B_{0}=1 ; \\
& a_{i}=i /\left(2 x^{2}\right), \quad b_{i}=1 \quad \text { for } i=0,1,2,3, \ldots ;  \tag{10}\\
& b_{0}+\frac{a_{1} \mid}{\mid b_{1}}+\frac{a_{2} \mid}{\mid b_{2}}+\cdots+\frac{a_{n} \mid}{\mid b_{n}}=\frac{A_{n}}{B_{n}} .
\end{align*}
$$

Then $A_{n}$ and $B_{n}$ are given recursively by the relations

$$
\left.\begin{array}{rl}
A_{n} & =b_{n} A_{n-1}+a_{n} A_{n-2}  \tag{11}\\
B_{n} & =b_{n} B_{n-1}+a_{n} B_{n-2}
\end{array}\right\}, \quad n=1,2,3, \ldots
$$

At smaller values of $x$, the convergence of (9) becomes slower. To overcome this, we have used double Aitken $\delta^{2}$ extrapolation (Burden et al. [11, pp. 56-57]) as follows. If $y_{j}$ is the approximation obtained by taking $j$ terms of (9), then the sequences of numbers

$$
\begin{align*}
& y_{j}^{\prime}=y_{j}-\left(y_{j}-y_{j-1}\right)^{2} /\left(y_{j}-2 y_{j-1}+y_{j-2}\right),  \tag{12}\\
& y_{j}^{\prime \prime}=y_{j}^{\prime}-\left(y_{j}^{\prime}-y_{j-1}^{\prime}\right)^{2} /\left(y_{j}^{\prime}-2 y_{j-1}^{\prime}+y_{j-2}^{\prime}\right),
\end{align*}
$$

converge progressively faster to (9). Use of (12) with (9)-(11) improved the fit of the Chebyshev approximation by about five orders of magnitude.

We have used (7) for $x \leqslant 2.83$ and (9)-(12) for $x>2.83$. This yielded values of $f(x)$ accurate to at least 23D for all $x$. The resulting Chebyshev coefficients,
generated by (3)-(5), are shown in Table 1. Use of these in (2) gives an approximation which reproduces $f(x)$ to at least 22D for all $x$. Schonfelder [5] generates Chebyshev coefficients using a different method [12], [13] in which an expression equivalent to (2) is substituted into a linear differential equation satisfied by the given function. Together with the boundary conditions satisfied by the same function, this procedure generates an infinite set of simultaneous linear equations for the $c_{j}$, a truncated version of which is then solved.

Table 1

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